

Elektrostatik:

Coulomb-Kraft:  $\vec{F} = \frac{q_1 q_2}{4\pi\epsilon_0 r^2} \cdot \frac{\vec{r}}{r} = q_1 \cdot \vec{E}$

Potential:  $\phi(r) = \frac{Q}{4\pi\epsilon_0 r} \Rightarrow \vec{E} = -\nabla\phi$

Superpositionsprinzip gilt für:  $\vec{F}, \vec{E}, \phi$

Spannung:  $U_{1,2} = \int_{\vec{r}_1}^{\vec{r}_2} \vec{E} \cdot d\vec{s} = \phi(\vec{r}_1) - \phi(\vec{r}_2)$

$E_{Pot} = \int_{\vec{r}_1}^{\vec{r}_2} \vec{F} \cdot d\vec{s} = qU$

(Raum-)Ladungsdichte  $\rho: \int_V \rho \cdot dV = Q$

El. Fluss:  $\phi_{el} = \int_A \vec{E} d\vec{a} = EA \quad \oint_A \vec{E} d\vec{A} = \frac{1}{\epsilon_0} Q$

Gaußscher Satz:  $\oint_A \vec{v} d\vec{a} = \int_V \vec{\nabla} v dV$

Poisson-Gleichung:  $-\nabla^2 \phi = \frac{\rho}{\epsilon_0}$

Im Leiter:  $\vec{E}_{\parallel} = 0, \vec{E}_{\perp} = \vec{E}, \vec{E}_{Innen} = 0$

Kondensator:  $C = \frac{Q}{U} \quad E_C = \frac{1}{2} QU \quad I = CU$

Platten:  $E = \frac{U}{d} = \frac{Q}{\epsilon_0 A} \quad C = \epsilon_0 \frac{A}{d}$

Energiedichte:  $\ell = \frac{1}{2} \epsilon_0 E^2 \quad E_C = \ell Ad = \frac{1}{2} CU^2$

Diel.konst.:  $Q = Q_0 + Q_{pol} = \frac{1}{\epsilon} Q_0 \quad C_{Diel} = \epsilon C_{Vak}$

$E = \frac{1}{\epsilon_0} (\underbrace{Q_0}_A + \underbrace{Q_{pol}}_{-P}) = \frac{Q_0}{\epsilon_0 A} \Rightarrow \vec{E} = \frac{1}{\epsilon_0} (\vec{D} - \vec{P})$

lineares Mat.:  $\vec{D} = \epsilon \epsilon_0 \vec{E} \quad \vec{P} = \alpha \vec{E} \quad \alpha = \epsilon - 1$

$\vec{\nabla} \vec{E} = \frac{\rho}{\epsilon_0} \quad \vec{\nabla} \vec{D} = \rho_0 \quad \vec{\nabla} \vec{P} = -\rho_{pol} \quad \rho = \rho_0 + \rho_{pol}$

Strom:

Stromstärke:  $I = \frac{dQ}{dt} = \int_A \vec{j} d\vec{a} = nqAv$

$\vec{j} = nq\vec{v} = \sigma \vec{E} \quad \vec{\nabla} \vec{j} = -\frac{\partial \rho}{\partial t}$

spez. Leitfähigkeit:  $\sigma = \frac{1}{\rho}$  spez. Widerstand:  $\rho$

Widerstand:  $R = \frac{U}{I}$  Leiter:  $R = \rho \frac{\ell}{A} \quad \rho = \frac{E}{j}$

Suszeptibilität:  $\vec{j} = \chi \vec{E}$

$R(t) = \int \chi(t-t') F(t') dt' \sim e^{-\frac{t}{\tau}}$

Kontinuitätsgleichung:  $\oint \vec{j} d\vec{a} = -\frac{d}{dt} \int_V \rho dV$

Leistung:  $P_{el} = \frac{dE_{el}}{dt} = \frac{dq}{dt} U = IU$

Kirchhoffsche Regeln - Maschen:  $U_0 = \sum_k U_k$

-Knoten:  $\sum I_k = 0$

Magnetismus:

$\oint_C \vec{B} d\vec{s} = \mu_0 I \neq 0 \quad \oint_{\partial A} \vec{B} d\vec{s} = \int_A \vec{\nabla} \times \vec{B} d\vec{a}$

$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j} = -\Delta \vec{A} \quad \phi_m = \int \vec{B} d\vec{a}$

$\vec{\nabla} \vec{B} = \vec{\nabla} \vec{A} = 0$

Vektorpotential:  $\vec{B} = \vec{\nabla} \times \vec{A}$

Biot-Savart:  $\vec{B} = \frac{\mu_0}{4\pi} \int_C I \frac{(\vec{r}' - \vec{r}) \times d\vec{s}}{|\vec{r}' - \vec{r}|^3}$

$= \frac{\mu_0}{4\pi} \int_V \frac{\vec{j} \times |\vec{r}' - \vec{r}|}{|\vec{r}' - \vec{r}|^3} dV$

Leiter außen:  $B(r) = \frac{\mu_0}{2\pi r} I$  Innen:  $B(r) = \frac{\mu_0}{2\pi R_0^2} r I$

Leiterschleife:  $B(z) = \frac{\mu_0 I R^2}{2\sqrt{R^2 + z^2}} \Rightarrow B_{Zent} = \frac{\mu_0 I}{2R}$

$B_{\infty} \approx \frac{\mu_0 I R^2}{2z^3} \quad m = \pi I R^2$

mag. Dipol:  $B(\vec{r}) = \frac{\mu_0}{4\pi} \frac{3\vec{r}(\vec{r} \cdot \vec{m}) - \vec{m} r^2}{r^5}$

mag. Moment:  $\vec{m} := I \vec{A}$  Zylinderspule:  $B = \mu_0 I \frac{N}{L}$

Lorentz-Kraft - Teilchen:  $\vec{F} = q(\vec{v} \times \vec{B})$

-Dünner Draht:  $F = IB_{\perp} \ell$

- Leiterstück:  $\vec{F} = \int_V (\vec{j} \times \vec{B}) dV$

- Zwischen zwei Leitern:  $F = \frac{\mu_0 I_1 I_2 L}{2\pi d}$

Massenspektrometer:  $F_{\omega} = m \frac{v^2}{R} \quad R = \frac{mv}{qB}$

Hall-Spannung:  $U_H = -\frac{(\vec{j} \times \vec{B}) \cdot d}{nq} = \frac{IB}{nqd}$

$\vec{\tau} = \vec{m} \times \vec{B} \quad E_{pot} = -\vec{m} \cdot \vec{B}$  (m, B konstant)

Faraday-Kraft:  $\vec{F} = \vec{\nabla}(\vec{m} \cdot \vec{B})$  im homogenen ext. B-Feld

Magnetisierung:  $\vec{M} = \frac{\vec{m}}{V} \quad \vec{\nabla} \times \vec{M} = \vec{j}_{Mag}$

Permeabilität:  $\vec{j} = \vec{j}_0 + \vec{j}_{Mag} = \mu \vec{j}_0$

$\vec{B} = \mu_0(\vec{H} + \vec{M}) = \mu \mu_0 \vec{H} \quad \vec{\nabla} \times \vec{H} = \vec{j}_0$

Zeitlich veränderliche Felder:

$U_{ind} = -\frac{d}{dt} \phi_m$  Selbstinduktivität L:  $U_{ind} = -L \frac{dI}{dt}$

$\phi_m = LI$

Zylinderspule:  $L = \mu_0 \frac{N^2 A}{\ell}$

$E_L = \frac{1}{2} LI^2$  Energiedichte:  $e_m = \frac{1}{2\mu_0} B^2$

Maxwell-Gleichungen:

$\vec{\nabla} \vec{E} = \frac{1}{\epsilon_0} \rho \quad \oint \vec{E} d\vec{a} = \frac{1}{\epsilon_0} \int \rho dV$

$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \oint \vec{E} d\vec{s} = -\frac{d}{dt} \int \vec{B} d\vec{a}$

$\vec{\nabla} \vec{B} = 0 \quad \oint \vec{B} d\vec{a} = 0 \quad \vec{\nabla} \times \vec{B} = \mu_0 \vec{j} + \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$

$\oint \vec{B} d\vec{s} = \mu_0 \int \vec{j} d\vec{a} + \frac{1}{c^2} \frac{d}{dt} \int \vec{E} d\vec{a} \quad c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$

Stromkreise

Reihe:  $R = \sum R_i \quad \frac{1}{C} = \sum \frac{1}{C_i} \quad L = \sum L_i$

Parallel:  $\frac{1}{R} = \sum \frac{1}{R_i} \quad C = \sum C_i \quad \frac{1}{L} = \sum \frac{1}{L_i}$

U-quelle:  $U_{aus} = U_0 - R_i I \quad I\text{-quelle: } I = I_0 - \frac{1}{R_i} U$

Kondensator:  $\tau = RC \quad I_0 = \frac{U_0}{R}$

$I_{an}(t) = I_0 e^{-\frac{t}{\tau}} \quad I_{aus}(t) = -I_0 e^{\frac{t}{\tau}}$

Spule:  $\tau = \frac{L}{R}$

$I_{an}(t) = \frac{U_0}{R} (1 - e^{-\frac{t}{\tau}}) \quad I_{aus}(t) = \frac{U_0}{R} e^{-\frac{t}{\tau}}$

Wechselstromkreise:  $j = -i$

$U(t) = U_0 \cos(\omega t + \varphi_U) \quad I(t) = I_0 \sin(\omega t + \varphi_I)$

$A \cos(\omega t + \varphi) = a' \cos(\omega t) + a'' \sin(\omega t) = \Re(\hat{A} e^{j\omega t})$

R:  $\varphi_U = \varphi_I \quad C: \varphi_U - \varphi_I = \frac{\pi}{2} \quad L: \varphi_U - \varphi_I = -\frac{\pi}{2}$

Scheinleistung:  $P(t) = U(t)I(t) \quad \bar{P} = \frac{1}{T} \int U(t)I(t) dt$

$\Rightarrow P(t) = U_0 I_0 \cos(\omega t) \underbrace{(\cos(\varphi_I) \cos(\omega t))}_{\text{Wirkleistung}} - \underbrace{(\sin(\varphi_I) \sin(\omega t))}_{\text{Blindleistung}}$

$\Rightarrow \bar{P} = \frac{U_0 I_0}{2} \cos(\varphi)$

$\hat{Z} = \frac{\hat{U}}{\hat{I}} \quad \hat{Z}_R = R \quad \hat{Z}_C = -\frac{j}{\omega C} \quad \hat{Z}_L = j\omega L$

Zweiter Ü-funktion:  $\hat{F}(\omega) = \frac{\hat{U}_{aus}}{\hat{U}_{ein}}$

Filter: Hochpass, Tiefpass, Bandpass, Allpass

Grenzfrequenz:  $A(\omega_g) = \frac{1}{2} A_0$

Elektromagnetische Wellen:

$f(x, t) = A \cos(\omega t - kx + \varphi)$

$\vec{f}(t) = \vec{A} \cos(\omega t - \vec{k} \cdot \vec{x}) = \hat{A} e^{i(\vec{k} \cdot \vec{x} - \omega t)}$

Energiedichte:  $e = \frac{1}{2} \epsilon_0 \vec{E}^2(t) + \frac{1}{2\mu_0} \vec{B}^2(t)$

$e = \epsilon_0 E_0^2 = \frac{1}{\mu_0} B_0^2$

Energiestromdichte (=Poynting-Vektor):

$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} = I \frac{\vec{k}}{k}$

Intensität:  $I = |\vec{S}| = ec = c \epsilon_0 E_0^2$

Strahlungsdruck:  $P = \frac{I}{c}$

Impulsdichte:  $\vec{\pi} = \epsilon_0(\vec{E} \times \vec{B}) = \frac{1}{c^2} \vec{S}$

Superpositionsprinzip gilt nicht für I

Periodische ebene Wellen:  $\omega = c|k| \Rightarrow k = \frac{2\pi}{\lambda}$

Doppelspalt:  $I_{max}$  für  $\delta = n\lambda \Rightarrow \delta = d \cos(\alpha)$

Bragg-Bedingung:  $2d \sin(\alpha) = m\lambda \quad m, n \in \mathbb{N}$

Metallreflexion:  $\vec{E}_{refl} = -\vec{E}_{ein} \quad \vec{B}_{refl} = \vec{B}_{ein}$

Lineare Polarisation:

$\vec{E} = E_0 \begin{pmatrix} \cos(\varphi) \\ \sin(\varphi) \\ 0 \end{pmatrix} \cos(\omega t - kz + \varphi)$

Zirkulare P.:  $\vec{E} = E_0 \begin{pmatrix} \cos(\omega t - kz + \varphi) \\ \cos(\omega t - kz + \varphi \pm \frac{\pi}{2}) \\ 0 \end{pmatrix}$

- Rechtsdrehend:  $\sigma^+$  Linksdrehend:  $\sigma^-$

Elliptische P.:  $\vec{E} = \begin{pmatrix} E_x \cos(\omega t - kz + \varphi) \\ E_y \cos(\omega t - kz + \varphi \pm \frac{\pi}{2}) \\ 0 \end{pmatrix}$

Fernfeldstrahlung Hertzscher Dipol:

$\vec{E} = \frac{\omega^2}{4\pi\epsilon_0 c^2 r} ((\vec{n} \times \vec{p}_0) \times \vec{n}) e^{i(kr - \omega t)}$

$\vec{B} = \frac{\omega^2}{4\pi\epsilon_0 c^3 r} (\vec{n} \times \vec{p}_0) e^{i(kr - \omega t)} \quad \vec{n} = \frac{\vec{r}}{r} = \frac{\vec{k}}{k}$

$\Rightarrow \vec{E} = c(\vec{B} \times \vec{n})$  (E und B sind in Phase)

$\vec{S} = \frac{1}{\epsilon_0 c^3} [\frac{p_0 \omega^2 \sin(\theta)}{4\pi r} \cos(\omega t - kr)]^2 \vec{n}$

$P = \frac{p_0^2 \omega^4}{12\pi \epsilon_0 c^3} \frac{dP}{d\Omega} \sim \sin^2(\theta)$

Strahlungsleistung beschl. Ladung:  $P = \frac{2q^2}{3\epsilon_0 c^3} \dot{v}^2$

In Materie: - Brechungsindex:  $n = \frac{c_0}{c} = \sqrt{\epsilon \mu} \quad c = \frac{1}{\sqrt{\epsilon \epsilon_0 \mu \mu_0}}$

n, ε, c frequenzabhängig

Optik: HIER KÖNNTE IHRE WERBUNG STEHEN!!

Konstanten/Einheiten:

$[Q] = C = As \quad [\vec{E}] = \frac{N}{As} = \frac{V}{m} \quad [U] = \frac{Nm}{As} = V$

$[C] = \frac{C}{V} = F \quad [\sigma] = \frac{A}{Vm} \quad [L] = \frac{Vs}{A} = H$

$[\vec{D}] = \frac{C}{m^2} \quad [J] = A = \frac{C}{s} \quad [R] = \Omega = \frac{V}{A}$

$[P] = VA = W \quad [M] = [H] = \frac{A}{m} \quad [B] = T = \frac{N}{C \frac{m}{s}} = \frac{N}{Am}$

$\epsilon_0 = 8,854 \cdot 10^{-12} \frac{A^2 s^4}{kg m^3} \quad e = 1,6021764 \cdot 10^{-19} C$

$k = \frac{1}{4\pi \epsilon_0} = 8,99 \cdot 10^9 \frac{Nm^2}{C^2} \quad \mu_0 = 4\pi \cdot 10^{-7} \frac{Vs}{m^2}$

Other shit:

$\text{grad } f = \vec{\nabla} f \quad \text{div } \vec{f} = \vec{\nabla} \cdot \vec{f} \quad \text{rot } \vec{f} = \vec{\nabla} \times \vec{f}$

Gaußscher Satz (Mathe):  $\int_{\Omega} (\vec{\nabla} \cdot \vec{F}) d\vec{r} = \oint_{\partial \Omega} \vec{F} d\vec{A}$

Stokes:  $\int_A (\vec{\nabla} \cdot \vec{F}) d\vec{A} = \oint_{\partial A} \vec{F} d\vec{l}$

Filter/2-Tore:

Übertragungsfunktion:  $\hat{F}(\omega) = \frac{\hat{U}_{aus}}{\hat{U}_{ein}}$

$\hat{I}_{ein,1} = \hat{I}_{ein,2}, \hat{I}_{aus,1} = \hat{I}_{aus,2}$

Amplitudengang:  $A(\omega) = |\hat{F}(\omega)|$

Phasengang:  $\varphi = \arctan(\frac{\Im \hat{F}}{\Re \hat{F}})$

Integralbla:  $dA = dx dy = r dr d\phi \quad dV = r^2 dr \sin(\theta) d\theta d\phi$

$\int_{\Omega} \vec{\nabla} \vec{F} dV = \oint_{\partial \Omega} \vec{F} d\vec{A} \quad \int_A \vec{\nabla} \times \vec{F} d\vec{A} = \oint_{\partial A} \vec{F} d\vec{l}$

$\int_{\gamma} \vec{F} d\vec{l}(t) = \int_{t_1}^{t_2} (\vec{F} \frac{d\vec{l}}{dt}) dt$

Kugelkoordinaten:

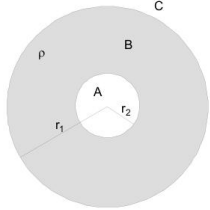
$x = r \sin(\theta) \cos(\phi) \quad \phi \in [0, 2\pi)$

$y = r \sin(\theta) \sin(\phi) \quad \theta \in [0, \pi)$

$z = r \cos(\theta)$

Tera	T	10 <sup>12</sup>	Giga	G	10 <sup>9</sup>
Mega	M	10 <sup>6</sup>	Kilo	k	10 <sup>3</sup>
Piko	p	10 <sup>-12</sup>	Nano	n	10 <sup>-9</sup>
Mikro	μ	10 <sup>-6</sup>	Milli	m	10 <sup>-3</sup>

Unendlicher homogener geladener Hohlzylinder, Flächenladungsdichte  $\rho$



E-Feld B:  $q = \rho V = \pi \rho \ell (r_2^2 - r_1^2) \quad 2\pi r \ell E = \frac{q}{\epsilon_0}$

E-Feld C:  $q = \rho V = \pi \rho \ell (r_1^2 - r_2^2)$

Zylinderkondensator der Länge L (Bild wie oben):

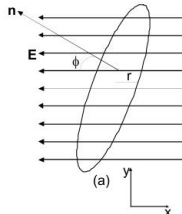
$\rho = \frac{Q}{V} \Rightarrow Q = 2\pi r_1 L \rho_1 \Rightarrow \rho_2 = \frac{r_1}{r_2} \rho_1$

$\oint \vec{E} d\vec{a} = \frac{Q}{\epsilon_0} \Rightarrow E 2\pi r L = \frac{\rho_2 2\pi r_2 L}{\epsilon_0} \Rightarrow E = \frac{\rho_1 r_1}{\epsilon_0 r}$

$U = \int \vec{E} d\vec{s} = \int_{r_2}^{r_1} \frac{\rho_2 r_2}{\epsilon_0 r} = \frac{\rho_2 r_2}{\epsilon_0} \ln\left(\frac{r_1}{r_2}\right)$

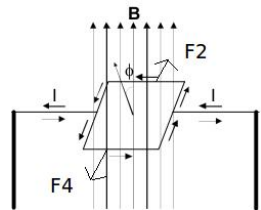
$C = \frac{Q}{U} = \frac{2\pi \epsilon_0 L}{\ln\left(\frac{r_1}{r_2}\right)}$

E-Fluss durch Kreisfläche:



$\phi_{el} = \int_A \vec{E} d\vec{a} = E \cos(\phi) \pi r^2$

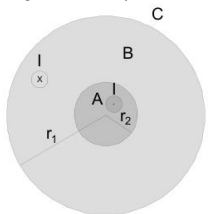
Leiterschleife im B-Feld:



$F = I \int_C \vec{B} \times \vec{s} = I \int_0^a B \sin(\phi) ds \quad (F_2, F_4 : 2\phi)$

$\tau_{F1} = \tau_{F3} = 0 \Rightarrow \tau = \vec{r}_2 \times \vec{F}_2 + \vec{r}_4 \times \vec{F}_4$

Magnetfeld eines zylindrischen Leiters:



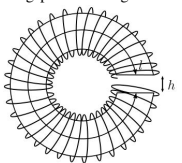
Hohl:  $B: \vec{E} = 0 \quad C: \oint \vec{B} d\vec{s} = B 2\pi r = \mu_0 I \Rightarrow B = \frac{\mu_0 I}{2\pi r}$

Mitte Hohl:  $A: B 2\pi r = \mu_0 \frac{I}{\pi r^2} \pi r^2 \Rightarrow B = \frac{\mu_0 I r}{2\pi r^2}$

$-B: B 2\pi r = \mu_0 I \quad C: B = 0$

Voll:  $A, C: s.o. B: B 2\pi r = \mu_0 \left(I - \frac{I}{\pi(r_1^2 - r_2^2)} \pi(r^2 - r_2^2)\right)$

Ringspule der Länge  $\ell$  mit Lücke  $h$  und Eisenkern



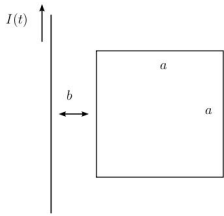
$\oint \vec{B} d\vec{A} = 0 \Rightarrow \mu_0 H_{gap} = \mu \mu_0 H_{Kern}$

$NI = \oint \vec{H} d\vec{r} = H_{Kern}(\ell - h) + H_{gap} h = H_{gap} \left(h + \frac{\ell - h}{\mu}\right)$

$\Rightarrow H_{gap} = \mu H_{Kern} = \mu \frac{NI}{\ell + h(\mu - 1)}$

$B_{gap} = \mu_0 H_{gap} \quad B_{Kern} = \mu \mu_0 H_{Kern}$   
 $M_{gap} = 0 \quad M_{Kern} = (\mu - 1) H_{Kern}$

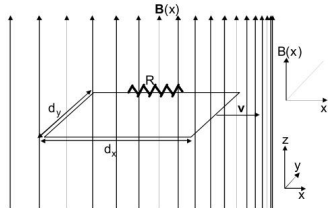
Induktion in quadratischer Leiterschleife, Draht:  $I(t) = I_0 \sin(\omega t)$



$U_{ind} = -\frac{d}{dt} \oint \vec{B} d\vec{A} = -\frac{\mu_0}{2\pi} \frac{d}{dt} I(t) \int_0^a \int_b^{b+a} \frac{1}{r} dr dz$   
 $= -\frac{\mu_0 I_0 a \omega}{2\pi} \ln\left(\frac{b+a}{b}\right) \cos(\omega t)$

$U_{ind} = -L \frac{dI}{dt} \Rightarrow L = \frac{\mu_0}{2\pi} a \ln\left(\frac{b+a}{b}\right)$

Bewegende Leiterschleife im Magnetfeld mit konstantem Gradient,  $B(x) = B_0 x$

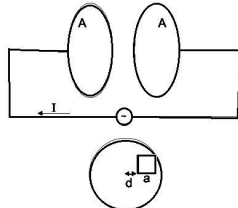


$\phi = \int_A \vec{B} d\vec{a} = \int_0^y \int_x^{x+dx} B(x') = B_0 dx dy \left(x + \frac{dx}{2}\right)$

$U = RI = \frac{d\phi}{dt} = B_0 dx dy v \Rightarrow I = \frac{B_0 dx dy v}{R}$

$P = I^2 R = Fv \Rightarrow F = \frac{B_0 dx dy^2 v}{R}$

Verschiebungsstrom im Kondensator mit  $I(t) = I_0 \cos(\omega t)$  und  $Q(0) = 0$

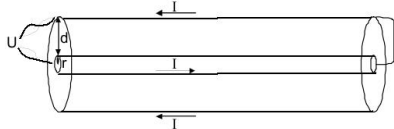


$I = \frac{dQ}{dt} \Rightarrow \int_0^t dq = \int_0^t I(t') dt' \Rightarrow Q(t) = \frac{I_0}{\omega} \sin(\omega t)$

$\oint \vec{E} d\vec{a} = E(t)A = \frac{Q(t)}{\epsilon_0} \Rightarrow E(t) = \frac{I_0 \sin(\omega t)}{\epsilon_0 \omega A}$

$I_D = \epsilon_0 \frac{\partial E(t)}{\partial t} = \frac{I_0}{A} \cos(\omega t)$

Selbstinduktion im Koaxialkabel

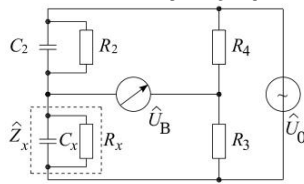


$\oint \vec{B} d\vec{a} = B 2\pi r' = \mu_0 I$

$\phi = \int_A \vec{B} d\vec{a} = \int_0^\ell dz \int_r^{r+d} dr' B = \frac{\mu_0 \ell}{2\pi} \ln\left(1 + \frac{d}{r}\right) I$

$U_{ind} = -\frac{d\phi}{dt} = -\frac{\mu_0 \ell}{2\pi} \ln\left(1 + \frac{d}{r}\right) \frac{dI}{dt}$

Wheatstone-Brücke mit Spannfrequenz  $\omega$



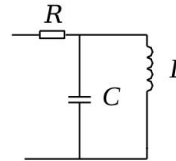
$R_x, R_2 \mapsto \infty, \hat{U}_B = 0 : \frac{\hat{Z}_x}{\hat{Z}_2} = \frac{\hat{Z}_3}{\hat{Z}_4} \quad \hat{Z}_3 = R_3, \hat{Z}_4 = R_4$

$\hat{Z}_x = \frac{1}{j\omega C_x}, \hat{Z}_2 = \frac{1}{j\omega C_2} \Rightarrow C_x = \frac{R_4}{R_3} C_2$

mit  $R_x, R_2: \frac{1}{\hat{Z}_x} = \frac{1}{R_x} + j\omega C_x \quad \frac{1}{\hat{Z}_2} = \frac{1}{R_2} + j\omega C_2$

Balanced:  $\frac{1}{R_x} + j\omega C_x = \frac{R_4}{R_3} \left(\frac{1}{R_2} + j\omega C_2\right)$

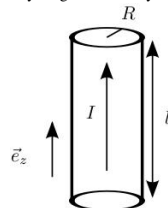
Impedanz bei  $U(t) = U_0 \cos(\omega t)$ :



$\hat{Z} = \hat{Z}_R + \hat{Z}_{LC} = \hat{Z}_R + \frac{1}{\frac{1}{\hat{Z}_L} + \frac{1}{\hat{Z}_C}} = R + j \frac{1}{\omega L - \frac{1}{\omega C}}$

$\bar{P} = \Re\left(\frac{1}{2} \hat{U} \hat{I}^*\right) = \frac{1}{2} U_0^2 \Re\left(\frac{1}{\hat{Z}}\right) = \frac{1}{2} U_0^2 \frac{R}{R^2 + \left(\frac{1}{\omega L} - \frac{1}{\omega C}\right)^2}$

Poynting Vektor im zylindrischen Leiter mit spez. Leitfähigkeit  $\sigma$ :



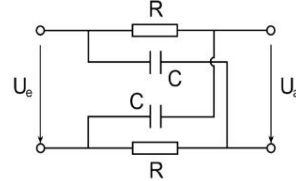
Innen:  $\vec{E} = \frac{U}{\ell} \vec{e}_z \quad \vec{B} = \frac{\mu_0 I r}{2\pi I R^2} \vec{e}_\phi \quad I = \sigma \vec{E} \vec{A} = \frac{\sigma U \pi R^2}{\ell}$

$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B} = -\frac{U^2 \sigma r}{2\ell^2} \vec{e}_r$

Oberfläche:  $r = R$  außen:  $\vec{S} = 0$

$P = \int \vec{S} d\vec{A} = S \int dA = S 2\pi R L$

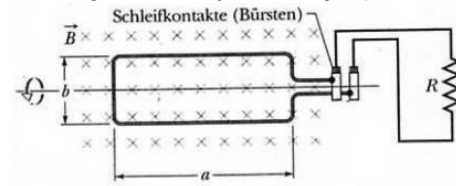
Filterkram:



$\hat{U}_{R2} = \hat{U}_e \frac{\hat{Z}_{R2}}{\hat{Z}_{C1} + \hat{Z}_{R2}} \quad \hat{U}_{C2} = \hat{U}_e \frac{\hat{Z}_{C2}}{\hat{Z}_{R1} + \hat{Z}_{C2}}$

$\hat{U}_a = \hat{U}_{C2} - \hat{U}_{R2} \quad \hat{F}(\omega) = \frac{\hat{U}_a}{\hat{U}_e} = \frac{-\frac{j}{\omega C}}{R - \frac{j}{\omega C}} - \frac{R}{R - \frac{j}{\omega C}}$   
 $= \frac{1 - j\omega CR}{1 + j\omega CR} \quad A(\omega) = |\hat{F}(\omega)| = 1$

Generator: Spule mit N Windungen und Drehfrequenz  $f$



$A = ab \sin(2\pi ft) \quad \phi_m = BNA \quad U(t) = -\frac{d\phi_m}{dt}$

EM-Dipol: Abstand  $r_1$ , Winkel  $\theta_1$  zur Dipolachse, misst  $E_0(r_1, \theta_1)$ . Bei  $r_2, \theta_2$ :

$E \sim \sin\left(\frac{\theta}{r}\right) \Rightarrow E_0(r_2, \theta_2) = E_0(r_1, \theta_1) \frac{r_1}{r_2} \sin(\theta_2)$

$e(r_2, \theta_2) = \epsilon_0 E^2(2) \quad \bar{e}(2) = \frac{1}{2} e(2)$

$|\vec{S}|(r_2, \theta_2, t) = c \epsilon_0 E^2(r_2, \theta_2, t) = c \epsilon_0 (E_0(2) \cos(\omega t + \varphi_0))^2$

$\bar{I}(2) = \bar{S} = \frac{1}{2} c \epsilon_0 E_0^2(2)$

...mit Frequenz  $\nu$ :

$E_0(r) \sim \frac{1}{r} \quad E_0 = c B_0 \Rightarrow B_0(r_2) = \frac{r_1 E_0(r_1)}{c r_2}$

$r_{max} = \frac{r_1 E_0(r_1)}{E_{0min}}$